# Mode bifurcation in the Rayleigh-Taylor instability of binary condensates

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# Normal modes of the interface Introduction Helmholtz equation in 2D(scaled units) Rayleigh-Taylor Instability(RTI) is an instability of an interface when a lighter fluid supports a heavier one in a gravitational field or -2 (2)some external potential. It can also occur when a lighter fluid pushes a heavier one. Occurs due to unfavourable energy conditions and as a result, the fluids tend to swap their positions. Leads to turbulent mixing of the two fluids as the perturbations at densities are low. the interface **grow exponentially**. the form of Mathieu equations **RTI in binary Bose-Einstein condensates Objectives:** ► We examine the generation and subsequent evolution of RTI in anisotropic two-species Bose-Einstein condensates(TBEC) in a pancake-shaped trap. Initiate instability, we tune intraspecies interaction between atoms. $q = \Lambda e \alpha/4, e = \sqrt{1 - 1/\alpha}$ ► We analytically study **normal modes** of interface using elliptic and lies in the domain $[0, 2\pi)$ . cylindrical coordinates. ► We find the normal modes undergo **bifurcation** at particular values **Allowed solutions** of anisotropy. and ratio of number of atoms. Phase separated pancake-shaped TBECs 15 2-D Gross-Pitaevskii(GP) equation(scaled units) $\left[-\nabla_{\perp}^{2}+V_{k}(x,y)+\sum_{i=1}^{\tilde{}}\mathcal{N}_{kj}|\phi_{j}(x,y)|^{2}\right]\phi_{k}(x,y)=i\frac{\partial\phi_{k}}{\partial t},$ $V(x, y) = \frac{m\omega^2}{2}(x^2 + \alpha^2 y^2)$ $\blacktriangleright$ In the initial state, at t = 0-5► We consider a system of <sup>85</sup>Rb–<sup>87</sup>Rb atoms ► $a_{11} = 460a_{B}, a_{22} = 99a_{B}, a_{12} = a_{21} = 214a_{B}$ ► $ce_2(v,q) = ce_2(v + \pi,q)$ and ► **ce**<sub>2</sub> ► Ce⁄ -10 0 10 -10 0 10 numerical results. x (units of a<sub>osc</sub>) **Growth Rate** of a<sub>o</sub> y (units 12 -12 0 x (units of $a_{osc}$ ) -12 0 12 **RTI** sets in when $n_1 > n_2$ .

$$(\nabla_{\perp}^2 + \Lambda_k^2)\phi_k = 0$$
  
 $\Lambda_k^2 = \tilde{\mu_k} - \alpha, \ \tilde{\mu_k} = \mu_k - \sum_{j=1}^2 \mathcal{N}_{kj} |\phi_j(x, y)|^2$ 

This equation is valid only at the interface or close to it, where

In elliptic cylindrical coordinates, Helmholtz equation assumes

$$\frac{d^2 U}{du^2} - \left[\mathcal{A} + \frac{1}{2}\Lambda^2 e^2 \alpha \left(1 - \cosh 2u\right)\right] U = 0,$$
$$\frac{d^2 \Theta}{dv^2} + \left[\mathcal{A} + \frac{1}{2}\Lambda^2 e^2 \alpha \left(1 - \cos 2v\right)\right] \Theta = 0.$$
$$a = \Lambda^2 e^2 \alpha / 4 \ e = \sqrt{1 - 1/\alpha^2}$$

 $\mathbf{v}$  represents the ellipse, v is the angular coordinate which varies

Solutions of the angular Mathieu equation, are the  $ce_{n-1}(v,q)$  and  $se_n(v, q)$  functions, cosine and sine elliptic functions, respectively.



The solution of interest is  $ce_2(v, q)$ , it satisfies

$$c_2(v,q) = \operatorname{ce}_2(-v,q)$$

$$p_2(v,q)$$
 is maximum at  $v=0$ 

• Undergoes a smooth **bifurcation** at higher values of q.

• At higher anisotropy parameter  $\alpha$ , instead of four there must be six mushroom shaped inward superfluid flow, which matches with our

$$s = \pm \left[\mathcal{A} + 2q \left(1 - \cos 2v\right)\right]^{\frac{1}{4}} \left[\frac{g(n_1 - n_2)}{n_1 + n_2}\right]^{\frac{1}{4}}$$

 $\mathbf{n}_1(n_2)$  refers to density of **outer(inner)** species.

**Numerical Results** 









# **Salient Features:**

- the trap.

### References



## Mode bifurcation and density profiles

Density profiles show early stages of RTI. Development of non-linear patterns on changing the anisotropy of

 $\blacktriangleright$  Bifurcation occurs at particular value of  $\alpha$ .

[1] Arko Roy, S. Gautam, D. Angom, arXiv: 1210.0381, (2012). [2] S. Gautam and D. Angom, Phys. Rev. A 81, 053616, (2010). [3] S. Gautam and D. Angom, J. Phys. B 43, 095302, (2010). [4] S. Gautam and D. Angom, J. Phys. B 44, 025302, (2011)

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